

Analysis of Unsteady Blood Flow Through a Stenosed Artery with Constant and Variable Viscosities

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ABSTRACT: *This paper presents a theoretical study of the analysis of unsteady blood flow with constant and variable viscosities through a stenosed artery using a third grade fluid model. Incorporated into the models are the slip velocity and externally applied magnetic field. The methods employed in solving the equations governing the unsteady blood flow models with constant and variable viscosities are the Galerkin's weighted residual and Forth order Runge-Kutta. Important flow parameters such as flow velocity, flow rate, shear stress and flow resistance have been computed. Graphical representation shows that, for both cases of unsteady blood flow models with constant viscosity and variable viscosity, magnetic field and shear thinning increases with flow resistance but decreases the flow velocity, flow rate and shear stress. Increases in slip velocity and shear thinning lead to increases in flow velocity, flow rate and shear stress but decrease the flow resistance. Other parameters that can positively influence the flow velocity are the pressure gradient and Reynold number. Finally, the velocity profile of unsteady blood flow model with constant viscosity is higher than that with variable viscosity.*

KEYWORDS: Unsteady blood flow, Variable viscosity, constant viscosity, magnetic field, slip velocity, stenosed artery, third grade fluid model and Galerkin's Weighted residual methods.

INTRODUCITON

In a recent time, blood flow through normal or stenosed arteries has gained serious attention of many researchers [1-5] because blood and blood vessels are substantial health risk factors in the development of many cardiovascular diseases. Blood flow in the human circulatory system depends upon the pumping action of the heart. That is, the heart moves blood efficiently through the branching networks of arteries, capillaries and veins and the lungs cycle here quite effectively through the branching pulmonary passage thereby keeping the cells of our bodies alive and functioning.

It is well known fact that at high shear rates and large diameter arteries, blood, being predominantly a suspension of erythrocytes in plasma, behave as a Newtonian fluid. Some of the researchers that investigated the flow of blood through stenosed arteries and treating blood

as Newtonian fluid are (Amit and Shrivastar [6], Tanwar and Varshney [7], Ellahi *et al.* [8]). Existing literature in this area also reveals that shear rate of blood is low in their stenosed region. A good many researchers (Hatami *et al* [9], Haleh *et al* [10], Aziz [11], Ikbal *et al* [12], Misra and Shit [13]) studied the non-Newtonian flow of blood from various perspectives.

In all the above investigations, only constant viscosity was considered. Variable viscosity of blood is another interesting area of study. The mathematical model of blood flow through a tapered artery with mild stenosis and variable viscosity was studied by Verma and Pariha [14]. Sanjeev and Chandraskhar [15] investigated hematocrit effect on the axisymmetric blood flow through stenosed arteries. Jimoh, *et al* [16] recently investigated hematocrit and slip velocity influence on third grade blood flow and heat transfer through a stenosed artery. They used weighted residual method to obtained the results in plotted graphs. The graphs reveals that, flow velocity and flow rate decreases while wall shear stress and resistance to flow increases when the hematocrit parameter increases.

The present study focused on the analysis of unsteady third grade blood flow with constant and variable viscosities through a stenosed artery. The externally applied magnetic field and slip velocity are also taken into consideration in this study.

MATHEMATICAL MODELS

The momentum equations describing the unsteady fluid flow models with constant viscosity and variable viscosity as obtained by Mohammed [17] and Jimoh [18] are respectively given as

$$\frac{\partial w}{\partial t} = \frac{\mu}{\rho r} \frac{\partial}{\partial r} \left(r \frac{\partial w}{\partial r} \right) + \frac{6\beta_3}{\rho} \left(\frac{\partial w}{\partial r} \right)^2 \frac{\partial^2 w}{\partial r^2} + \frac{2\beta_3}{\rho r} \left(\frac{\partial w}{\partial r} \right)^2 + \frac{\alpha_1}{\rho r} \frac{\partial^2 w}{\partial r \partial t} + \frac{\alpha_1}{\rho} \frac{\partial^3 w}{\partial r^2 \partial t} - \frac{\partial \hat{P}}{\rho \partial z} - \frac{\sigma \beta_0^2}{\rho} w \quad (2.1)$$

and

$$\begin{aligned} \frac{\partial w}{\partial t} = \frac{\mu_0}{\rho} \left[1 + N_1 \left(1 - \left(\frac{r}{R_0} \right)^m \right) \right] & \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial w}{\partial r} \right) + \frac{6\beta_3}{\rho} \left(\frac{\partial w}{\partial r} \right)^2 \left(\frac{\partial^2 w}{\partial r^2} \right) + \frac{2\beta_3}{r\rho} \left(\frac{\partial w}{\partial r} \right)^3 + \frac{\alpha_1}{r\rho} \frac{\partial^2 w}{\partial r \partial t} \\ & + \frac{\alpha_1}{\rho} \frac{\partial^3 w}{\partial r^2 \partial t} - \frac{1}{\rho} \frac{\partial \hat{P}}{\partial z} - \sigma \frac{\beta_0^2 w}{\rho} \end{aligned} \quad (2.2)$$

As a result of the stenosed artery as shown in figure 1, the corresponding slip conditions to (2.1) and (2.2) are respectively given as

$$w = w_s \quad \text{at} \quad r = R(z) \quad (2.3)$$

$$\frac{\partial w}{\partial r} = 0 \quad \text{at} \quad r = 0$$

and

$$\left. \begin{aligned} w &= w_{N_1 s} \quad \text{at} \quad r = R(z) \\ \frac{\partial w}{\partial r} &= 0 \quad \text{at} \quad r = 0 \end{aligned} \right\} \quad (2.4)$$

To non-dimensionalize equations (2.1) - (2.4), one introduces the following parameters and variables

$$\left. \begin{aligned} \bar{w} &= \frac{w}{d/t_0}, \quad y = r/R_0 \\ V_{01} &= \frac{w_s t_0}{d}, \quad V_{0N_1} = \frac{w_{N_1 s} t_0}{d} \end{aligned} \right\} \quad (2.5)$$

When equation (2.5) is substituted into (2.1) and (2.2), after simplifying one obtain respectively

$$\begin{aligned} \frac{\partial \bar{w}}{\partial \bar{t}} &= \frac{1}{RE_1} \cdot \frac{\partial}{\partial y} \left(y \frac{\partial \bar{w}}{\partial y} \right) + \Omega \left(6 \left(\frac{\partial \bar{w}}{\partial y} \right)^2 \frac{\partial^2 \bar{w}}{\partial y^2} + \frac{2}{y} \left(\frac{\partial \bar{w}}{\partial y} \right)^3 \right) + \Omega_1 \left(\frac{1}{y} \frac{\partial^2 \bar{w}}{\partial y \partial \bar{t}} + \frac{\partial^3 \bar{w}}{\partial y^2 \partial \bar{t}} \right) + G_1 - \\ &M_1 \bar{w} \end{aligned} \quad (2.6)$$

and

$$\begin{aligned} \frac{\partial \bar{w}}{\partial \bar{t}} &= \frac{1}{RE_{1N}} [1 + N_1(1 - y^m)] \cdot \frac{1}{y} \frac{\partial}{\partial y} \left(y \frac{\partial \bar{w}}{\partial y} \right) + \Omega_N \left(6 \left(\frac{\partial \bar{w}}{\partial y} \right)^2 \left(\frac{\partial^2 \bar{w}}{\partial y^2} \right) + \frac{2}{y} \left(\frac{\partial \bar{w}}{\partial y} \right)^3 \right) \\ &+ \Omega_{1N} \left(\frac{1}{y} \frac{\partial^2 \bar{w}}{\partial y \partial \bar{t}} + \frac{\partial^3 \bar{w}}{\partial y^2 \partial \bar{t}} \right) + G_{1N} - M_{1N} \bar{w} \end{aligned} \quad (2.7)$$

as the dimensionless momentum equations describing the unsteady blood flow models with constant and variable viscosities.

The dimensionless slip conditions corresponding to (2.6) and (2.7) can be simplified respectively as

$$\left. \begin{aligned} \bar{w} &= V_{01} & \text{at } y &= \frac{R(z)}{R_0} = R_b \\ \frac{\partial \bar{w}}{\partial y} &= 0 & \text{at } y &= 0 \end{aligned} \right\} \quad (2.8)$$

A

$$\left. \begin{aligned} \bar{w} &= V_{0N_1} & \text{at } y &= R_b \\ \frac{\partial \bar{w}}{\partial y} &= 0 & \text{at } y &= \end{aligned} \right\} \quad (2.9)$$

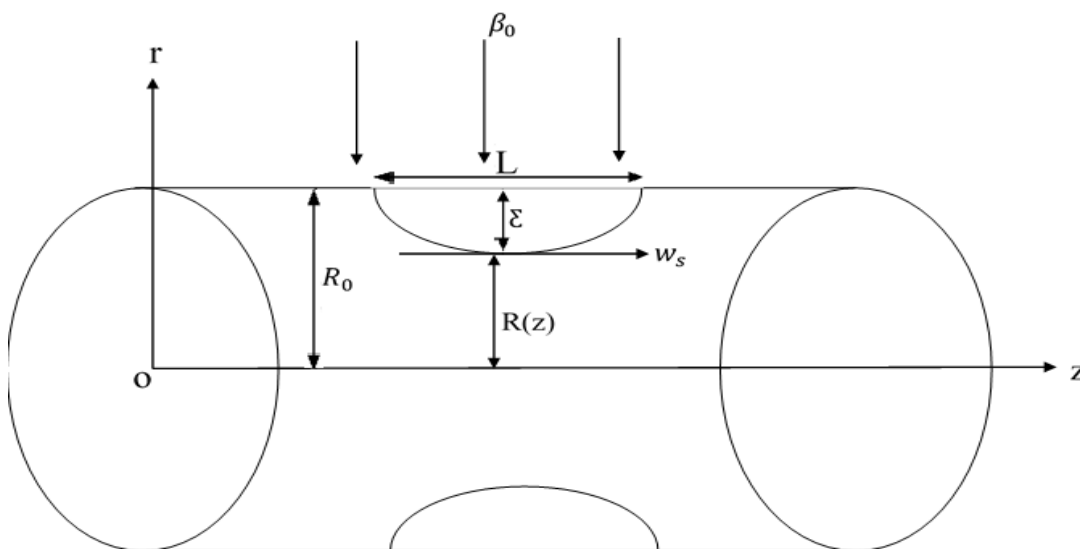


Figure1.Geometry of the stenosis

and has been described by Young [19] and Biswas [20]

$$\left. \begin{aligned} \frac{R(z)}{R_0} &= 1 - \frac{\varepsilon}{2R_0} \left[1 + \frac{\cos \pi z}{L} \right], \text{ for } |z| \leq L \\ R_0, & \text{ for } |z| > L \end{aligned} \right\} \quad (2.9)$$

METHODS OF SOLUTION

In order to solve (2.6) using Galerkin weighted residual method, one assume a solution of the form

$$\bar{w}(y, t) = a_0(t) + a_1(t)y + a_2(t)y^2 \quad (3.1)$$

Subjecting (3.1) to the slip conditions (2.8) and after simplifying yields

$$\bar{w}(y, t) = \frac{V_{01}y^2}{Rb^2} + a_0(t) \left(1 - \frac{y^2}{Rb^2} \right) + a_2(t) \frac{y^2}{Rb^2} \left(1 - \frac{y^2}{Rb^2} \right) \quad (3.2)$$

Using $\bar{r} = \frac{y}{Rb}$ in (3.2) and after simplifying and dropping the bar one obtain

$$w(r, t) = V_{01}r^2 + a_0(t)(1 - r^2) + a_2(t)r^2(1 - r^2) \quad (3.3)$$

The residual for equation (2.6) can be written as

$$\begin{aligned} RR_1(r, a_0(t), a_2(t)) &= \frac{\partial w}{\partial t} - G_1 - \frac{1}{RE_1} \frac{1}{r} \cdot \frac{\partial}{\partial r} \left(r \frac{\partial w}{\partial r} \right) - \Omega \left(6 \left(\frac{\partial w}{\partial r} \right)^2 \frac{\partial^2 w}{\partial r^2} + \frac{2}{r} \left(\frac{\partial w}{\partial r} \right)^3 \right) - \\ &\Omega_1 \left(\frac{1}{r} \frac{\partial^2 w}{\partial r \partial t} + \frac{\partial^3 w}{\partial r^2 \partial t} \right) + M_1 w \end{aligned} \quad (3.4)$$

substituting (3.3) into (3.4) and simplifying to obtain

$$\begin{aligned} RR_1(r, a_0(t), a_2(t)) &= -G_1 - \frac{1}{RE_1} (4V_{01} + 4a_2(t) - 4a_0(t) - 16a_2(t)r^2) + \\ &\dot{a}_0(t)(1 - r^2) + \dot{a}_2(t)r^2(1 - r^2) + 4\Omega_1(\dot{a}_0(t) - \dot{a}_2(t) + 4r^2\dot{a}_2(t)) - \Omega(48V_{01}^2r^2 - \end{aligned}$$

$$\begin{aligned}
& 144V_{01}^2a_0(t)r^2 + 144V_{01}^2a_2(t)r^2 - 480V_{01}^2a_2(t)r^4 + 144V_{01}a_0^2(t)r^2 - \\
& 40V_{01}a_0(t)a_2(t)r^2 + 960V_{01}a_0(t)a_2(t)r^4 + 144V_{01}a_2^2(t)r^2 - \\
& 144V_{01}a_2^2(t)r^4 + 4V_{01}a_2^2(t)r^6 + 144a_0^2a_2(t)r^2 - 480a_0^2(t)a_2(t)r^4 - 144a_0(t)a_2^2(t)r^2 + \\
& 960a_0(t)a_2^2(t)r^4 - 1344a_0(t)a_2^2(t)r^6 - 480a_2^3(t)r^4 + 1344a_2^3(t)r^6 - 48a_0^3(t)r^2 + \\
& 48a_2^3(t)r^2 - 1152a_2^3(t)r^8 - 128a_2^3(t)r^8 + 192V_{01}a_2^2(t)r^6 - 192a_0(t)a_2^2(t)r^6 + \\
& 192a_2^3(t)r^6 - 96V_{01}^2a_2(t)r^4 + 192V_{01}a_0(t)a_2(t)r^4 - 192V_{01}a_2^2(t)r^4 - \\
& 96a_0^2(t)a_2(t)r^4 + 192a_0(t)a_2^2(t)r^4 - 96a_2^3(t)r^4 + 16V_{01}^3r^2 - 48V_{01}^2a_0(t)r^2 + \\
& 48V_{01}a_0^2(t)r^2 - 96V_{01}a_0(t)a_2(t)r^2 + 48V_{01}a_2^2(t)r^2 - 16a_0^3(t) + 48a_0^2(t)a_2(t)r^2 - \\
& 48a_0(t)a_2^2(t)r^2 + 16a_2^3(t)r^2) - M_1(V_{01}r^2 + a_0(t) - a_0(t)r^2 + a_2(t)r^2 - a_2(t)r^4)
\end{aligned} \tag{3.5}$$

Differentiating (3.3) with respect to $a_0(t)$ and $a_2(t)$, one obtain $(1 - r^2)$ and $r^2(1 - r^2)$ respectively as the weight functions.

Taking the orthogonality of the residue $RR_1(a_0(t), a_2(t), r)$ with respect to the weight functions $(1 - r^2)$ and $r^2(1 - r^2)$, one respectively obtain the following systems of nonlinear first order differential equations;

$$\begin{aligned}
& (1848 + 9240\Omega_1)\dot{a}_0(t) + (264 - 1848\Omega_1)\dot{a}_2(t) + 29568\Omega a_0^3(t) + 23600\Omega a_2^3(t) - \\
& 88704\Omega V_{01}a_0^2(t) - 191136\Omega V_{01}a_2^2(t) + 29568\Omega a_0(t)a_0^2(t) + 25344\Omega a_2(t)a_0^2(t) - \\
& 146784\Omega V_{01}a_0(t)a_0^2(t) - \left(1848M_1 + 354816\Omega V_{01}^2 - \frac{9240}{RE_1}\right)a_0(t) - \left(264M_1 - \right. \\
& \left. 25344\Omega V_{01}^2 - \frac{1848}{RE_1}\right)a_2(t) = 2310G_1 + \frac{9240V_{01}}{RE_1} + 22176\Omega V_{01}^2 + 462M_1V_{01}
\end{aligned} \tag{3.6}$$

and

$$\begin{aligned}
 & (3432 + 24024\Omega_1)\dot{a}_0(t) + (1144 + 17160\Omega_1)\dot{a}_2(t) + 164736\Omega a_0^3(t) + \\
 & 4872384\Omega a_2^3(t) - 494208\Omega V_{01}a_0^2(t) - 690768\Omega V_{01}a_0^2(t) - 6877728\Omega a_0^2(t)a_2(t) - \\
 & 1677312\Omega a_2^2(t)a_0(t) - 1194336\Omega V_{01}a_0(t)a_2(t) - \left(3432M_1 - 1111968\Omega V_{01}^2 + \right. \\
 & \left. \frac{24024}{RE_1}\right)a_0(t) - \left(429M_1 - 329472\Omega V_{01}^2 + \frac{17160}{RE_1}\right)a_2(t) = 6006G_1 + \frac{24024}{RE_1} + 123552\Omega V_{01}^2 + \\
 & 2574M_1V_{01}
 \end{aligned} \tag{3.7}$$

By substituting the appropriate values of the parameters G_1 , V_{01} , RE_1 , Ω , Ω_1 , M_1 and t into (3.6) and (3.7) and solve using forth order Runge-Kutta method, one obtain the values for $a_0(t)$ and $a_2(t)$ which when substituted into (3.3), the velocity profiles were obtained which are shown in table 1.

Similarly, to solve equation (2.7) following the same procedure as indicated above, one can write residue for (2.7) as

$$\begin{aligned}
 RR_2(a_0(t), a_2(t), r) &= \frac{\partial w}{\partial t} - G_{1N} - \frac{1}{RE_{1N}}[1 + N_1(1 - r^m)]. \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial w}{\partial r} \right) - \\
 & \Omega_N \left(6 \left(\frac{\partial w}{\partial r} \right)^2 \left(\frac{\partial^2 w}{\partial r^2} \right) - \frac{2}{r} \left(\frac{\partial w}{\partial r} \right)^3 \right) - \Omega_{1N} \left(\frac{1}{r} \frac{\partial^2 w}{\partial r \partial t} + \frac{\partial^3 w}{\partial r^2 \partial t} \right) + M_{1N} w
 \end{aligned} \tag{3.8}$$

Taking the shape of the profile ($m = 2$) and substituting (3.3) into (3.8) and after simplifying in full to obtain.

$$\begin{aligned}
 RR_2(a_0(t), a_2(t), r) &= -64\Omega_N r^2 V_{01N}^3 + M_{1N} r^2 V_{01N}^3 + 1280a_2^3(t)\Omega_N r^8 - \\
 & 1536a_2^3(t)\Omega_N r^6 + 576a_2^3(t)\Omega_N r^4 - M_{1N} a_2(t)r^4 + 64a_0^3(t)\Omega_N r^2 - 64a_2^3(t)\Omega_N r^2 -
 \end{aligned}$$

$$\begin{aligned}
& M_{1N}a_0(t)r^2 + M_{1N}a_2(t)r^2 + 16\dot{a}_2(t)\Omega_{1N}r^2 + 4a_0(t) - 4a_2(t) + 16a_2(t)r^2 - 4V_{01N} - \\
& 1152a_0(t)a_2(t)\Omega_Nr^4V_{01N} + 384a_0(t)a_2(t)\Omega_Nr^2V_{01N} + 4N_1r^2V_{01N} - 16N_1a_2(t)r^4 - \\
& \dot{a}_2(t)r^4 - 4N_1a_0(t)r^2 + 20N_1a_2(t)r^2 - \dot{a}_0(t)r^2 + \dot{a}_2(t)r^2 + M_{1N}a_0(t) + 4\dot{a}_0(t)\Omega_{1N} - \\
& 4\dot{a}_2(t)\Omega_{1N} - G_{1N} - 4N_1V_{01N} + 4N_1a_0(t) - 4N_1a_2(t) + \dot{a}_0(t) + 1536a_0(t)a_2^2(t)\Omega_Nr^6 - \\
& 1536a_2^2(t)\Omega_Nr^6V_{01N} + 576a_0^2(t)a_2(t)\Omega_Nr^4 - 1152a_0(t)a_2^2(t)\Omega_Nr^4 + \\
& 1152a_2^2(t)\Omega_Nr^4V_{01N} + 576a_2(t)\Omega_Nr^4V_{01N}^2 - 192a_0^2(t)a_2(t)\Omega_Nr^2 - \\
& 192a_0^2(t)\Omega_Nr^2V_{01N} + 192a_0(t)a_2^2(t)\Omega_Nr^2 + 192a_0(t)\Omega_Nr^2V_{01N}^2 - 192a_2^2(t)\Omega_Nr^2V_{01N} - \\
& 192a_2(t)\Omega_Nr^2V_{01N}^2
\end{aligned} \tag{3.9}$$

where . is the differentiation with respect to time t.

The following systems of nonlinear first order differential equations were obtained by following the same procedures as indicated earlier.

$$\begin{aligned}
& 14784a_0^3(t)\Omega_NRE_{1N} + 12672a_0^2(t)a_2(t)\Omega_NRE_{1N} - 44352a_0^2(t)\Omega_NRE_{1N}V_{01N} + \\
& 14784a_0(t)a_2^2(t)\Omega_NRE_{1N} - 25344a_0(t)a_2(t)\Omega_NRE_{1N}V_{01N} + 44353a_0(t)\Omega_NRE_{1N}V_{01N}^2 + \\
& 2560a_2^3(t)\Omega_NRE_{1N} - 14784a_2^2(t)\Omega_NRE_{1N}V_{01N} + 12672a_2(t)\Omega_NRE_{1N}V_{01N}^2 - \\
& 14784\Omega_NRE_{1N}V_{01N}^3 + 924M_{1N}a_0(t)RE_{1N} + 132M_{1N}a_2(t)RE_{1N} + 231M_{1N}RE_{1N}V_{01N} + \\
& 4620\dot{a}_0(t)\Omega_{1N}RE_{1N} - 924\dot{a}_2(t)\Omega_{1N}RE_{1N} - 1155G_{1N}RE_{1N} + 3696N_1a_0(t) - \\
& 1584N_1a_2(t) - 3696N_1V_{01N} + 924\dot{a}_0(t)RE_{1N} + 132\dot{a}_2(t)RE_{1N} + 4620a_0(t) - \\
& 924a_2(t) - 4620V_{01N} = 0
\end{aligned} \tag{3.10}$$

and

$$\begin{aligned}
& -82368a_0^3(t)\Omega_N RE_{1N} - 164736a_0^2(t)a_2(t)\Omega_N RE_{1N} + 257104a_0^2(t)\Omega_N RE_{1N}V_{01N} - \\
& 122304a_0(t)a_2^2(t)\Omega_N RE_{1N} + 329472a_0(t)a_2(t)\Omega_N RE_{1N}V_{01N} - \\
& 247104a_0(t)\Omega_N RE_{1N}V_{01N}^2 - 33792a_2^3(t)\Omega_N RE_{1N} + 122304a_2^2(t)\Omega_N RE_{1N}V_{01N} - \\
& 164736a_2(t)\Omega_N RE_{1N}V_{01N}^2 + 82368\Omega_N RE_{1N}V_{01N}^3 - 1716M_{1N}a_0(t)RE_{1N} - \\
& 572M_{1N}a_2(t)RE_{1N} - 1287M_{1N}RE_{1N}V_{01N} - 12012\dot{a}_0(t)\Omega_{1N}RE_{1N} - \\
& 8580\dot{a}_2(t)\Omega_{1N}RE_{1N} + 3003G_{1N}RE_{1N} - 6864N_1a_0(t) - 2288N_1a_2(t) + 6864N_1V_{01N} - \\
& 1716\dot{a}_0(t)RE_{1N} - 572\dot{a}_2(t)RE_{1N} - 12012a_0(t) - 8580a_2(t) + 12012V_{01N} = 0
\end{aligned}
\tag{3.11}$$

Substituting the appropriate values of the parameters Ω_N , RE_{1N} , M_{1N} , Ω_{1N} , G_{1N} , V_{01N} , t and N_1 into (3.10) and (3.11) and solve using fourth order Runge-Kutta, one obtain the values for $a_0(t)$ and $a_2(t)$ which when substituted into (3.3) and simplified, one obtain the velocity profiles $w(r, t)$ which are shown in table 2.

Volume Flow Rate

The volume flow rate denoted by Q can be simplified as

$$Q = 12 \left[3V_{01}(R(z))^4 + a_0(t) \left(6(R(z))^2 - 3(R(z))^4 \right) + a_2(t) \left(3(R(z))^4 - 2(R(z))^6 \right) \right]
\tag{3.12}$$

Shear Stress

The shear stress denoted by τ_s can be simplified as

$$\tau_s = 2\mu R(Z) \left(V_{01} - a_0(t) + a_2(t) - 2R(Z)^2 a_2(t) \right) + 16R(Z)\beta_3 \left(V_{01} - a_0(t) + a_2(t) - 2(R(Z))^2 a_2(t) \right) \quad (3.13)$$

Resistance to Flow

The resistance to flow can be denoted as ψ can be simplified as

$$\psi = \frac{-\frac{\partial \hat{P}}{\partial z}}{12 \left[3V_{01}(R(z))^4 + a_0(t) \left(6(R(z))^2 - 3(R(z))^4 \right) + a_2(t) \left(3(R(z))^4 - 2(R(z))^6 \right) \right]} \quad (3.14)$$

Table 1: Values of the Parameters Used in the Numerical Results and the corresponding Velocity Profile for the Unsteady Blood Flow Model with constant viscosity.

Figures	G_1	V_{01}	RE_1	Ω	Ω_1	M_1	t	w(r, t)
2a	1.5	0.25	0.9	10	1	0.35	0.5	$-0.1437r^2 + 0.3437 + 0.0335r^2(1-r^2)$
	2.0	0.25	0.9	10	1	0.35	0.5	$-0.1858r^2 + 0.3858 + 0.0296r^2(1-r^2)$
	2.5	0.25	0.9	10	1	0.35	0.5	$-0.2290r^2 + 0.4290 + 0.0262r^2(1-r^2)$
3a	1.5	0.25	0.9	10	1	0.35	0.5	$-0.2261r^2 + 0.4261 + 0.0260r^2(1-r^2)$
	1.5	0.25	0.9	10	1	0.65	0.5	$-0.2174r^2 + 0.4174 + 0.0249r^2(1-r^2)$
	1.5	0.25	0.9	10	1	0.95	0.5	$-0.2090r^2 + 0.4090 + 0.0237r^2(1-r^2)$
4a	1.5	0.25	0.9	10	1	0.35	0.5	$-0.0572r^2 + 0.2572 + 0.0009r^2(1-r^2)$
	1.5	0.35	0.9	10	1	0.35	0.5	$-0.1538r^2 + 0.3538 + 0.0019r^2(1-r^2)$
	1.5	0.45	0.9	10	1	0.35	0.5	$-0.2510r^2 + 0.4510 + 0.0030r^2(1-r^2)$
5a	1.5	0.25	0.9	10	1	0.35	0.5	$-0.3246r^2 + 0.5246 + 0.0499r^2(1-r^2)$
	1.5	0.25	0.9	20	1	0.35	0.5	$-0.3108r^2 + 0.5108 + 0.0538r^2(1-r^2)$
	1.5	0.25	0.9	30	1	0.35	0.5	$-0.2848r^2 + 0.4848 + 0.0554r^2(1-r^2)$

	1.5	0.25	0.9	10	1	0.35	0.5	$-0.1246r^2 + 0.3246 + 0.0168r^2(1-r^2)$
6a	1.5	0.25	0.9	10	5	0.35	0.5	$-0.2016r^2 + 0.4016 + 0.0073r^2(1-r^2)$
	1.5	0.25	0.9	10	9	0.35	0.5	$-0.2542r^2 + 0.4542 + 0.0001r^2(1-r^2)$
	1.5	0.25	0.3	10	1	0.35	0.5	$-0.0053r^2 + 0.2053 + 0.0001r^2(1-r^2)$
7a	1.5	0.25	0.6	10	1	0.35	0.5	$-0.0094r^2 + 0.2094 + 0.0001r^2(1-r^2)$
	1.5	0.25	0.9	10	1	0.35	0.5	$-0.0091r^2 + 0.2091 + 0.0005r^2(1-r^2)$

Table 2: Values of the Parameters Used in the Numerical Results and the corresponding Velocity Profile for the Unsteady Blood Flow with Variable viscosity.

Figures	G_{1N}	V_{01N}	RE_{1N}	Ω_N	Ω_{1N}	M_{1N}	N_1	t	w(r, t)
	1.5	0.25	0.9	10	1	0.35	2	0.5	$0.2250 - 0.0250r^2 - 0.0119r^2(1-r^2)$
2b	2.0	0.25	0.9	10	1	0.35	2	0.5	$0.2588 - 0.0588r^2 - 0.0135r^2(1-r^2)$
	2.5	0.25	0.9	10	1	0.35	2	0.5	$0.2919 - 0.0919r^2 - 0.0138r^2(1-r^2)$
	1.5	0.25	0.9	10	1	0.35	2	0.5	$0.0218 - 0.0618r^2 - 0.0088r^2(1-r^2)$
3b	1.5	0.25	0.9	10	1	0.65	2	0.5	$0.0256 - 0.0540r^2 - 0.0086r^2(1-r^2)$
	1.5	0.25	0.9	10	1	0.95	2	0.5	$0.2503 - 0.0503r^2 - 0.0083r^2(1-r^2)$
	1.5	0.25	0.9	10	1	0.35	2	0.5	$0.3055 - 0.0555r^2 - 0.0111r^2(1-r^2)$
4b	1.5	0.35	0.9	10	1	0.35	2	0.5	$0.3954 - 0.0454r^2 - 0.0162r^2(1-r^2)$
	1.5	0.45	0.9	10	1	0.35	2	0.5	$0.4877 - 0.3773r^2 - 0.0206r^2(1-r^2)$
	1.5	0.25	0.9	10	1	0.35	2	0.5	$0.2974 - 0.0974r^2 - 0.0025r^2(1-r^2)$
5b	1.5	0.25	0.9	20	1	0.35	2	0.5	$0.2630 - 0.0630r^2 - 0.0082r^2(1-r^2)$

	1.5	0.25	0.9	30	1	0.35	2	0.5	$0.2582 - 0.0582r^2 - 0.0063r^2(1-r^2)$
	1.5	0.25	0.9	10	1	0.35	2	0.5	$0.2974 - 0.0974r^2 - 0.0025r^2(1-r^2)$
6b	1.5	0.25	0.9	10	5	0.35	2	0.5	$0.2872 - 0.0872r^2 - 0.0050r^2(1-r^2)$
	1.5	0.25	0.9	10	9	0.35	2	0.5	$0.2618 - 0.0618r^2 - 0.0088r^2(1-r^2)$
	1.5	0.25	0.3	10	1	0.35	2	0.5	$0.2381 - 0.0381r^2 - 0.0053r^2(1-r^2)$
7b	1.5	0.25	0.6	10	1	0.35	2	0.5	$0.2568 - 0.0568r^2 - 0.0067r^2(1-r^2)$
	1.5	0.25	0.9	10	1	0.35	2	0.5	$0.2618 - 0.0618r^2 - 0.0088r^2(1-r^2)$

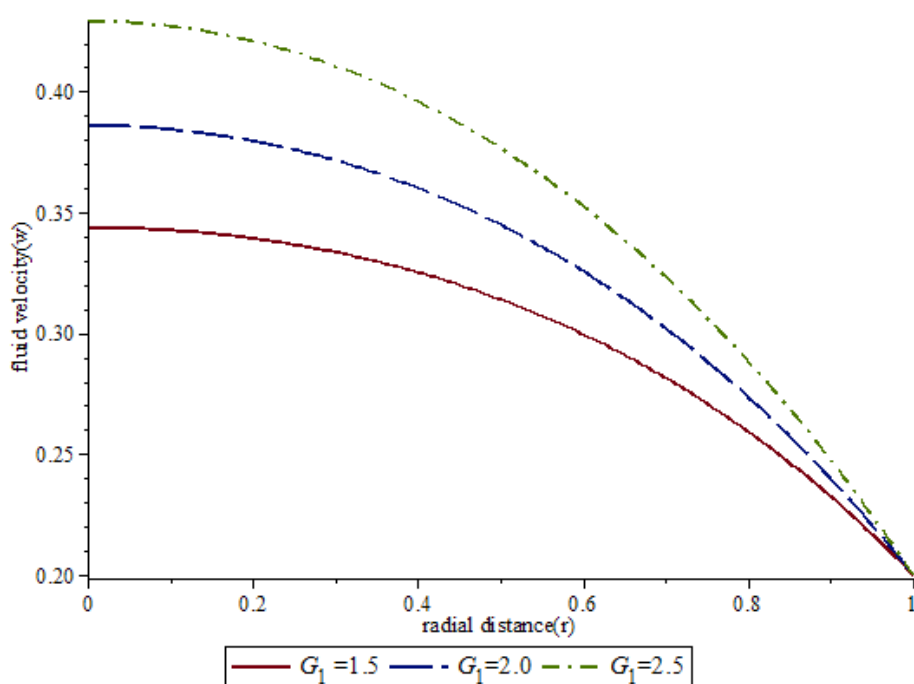


Figure 2a: Variation of Velocity Profile of the Unsteady Blood Flow Model with constant viscosity with increasing values of the Pressure Gradient in the radial direction.

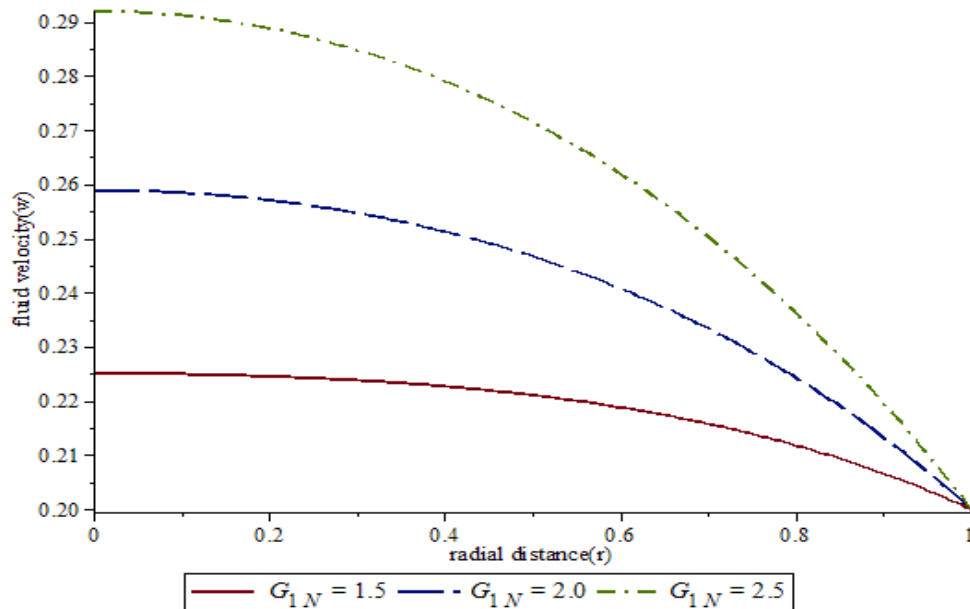


Figure 2b: Variation of Velocity Profile of the Unsteady Blood Flow Model with variable viscosity for various values of the Pressure Gradient in the radial direction.

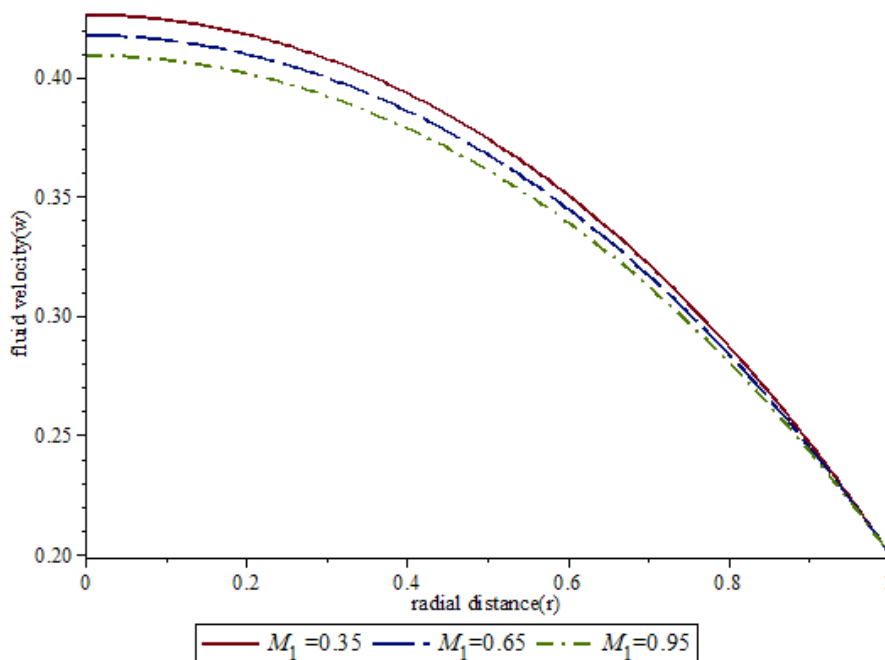


Figure 3a: Variation of Velocity Profile of the Unsteady Blood Flow Model with constant viscosity for various values of the Magnetic Field Parameter in the radial direction.

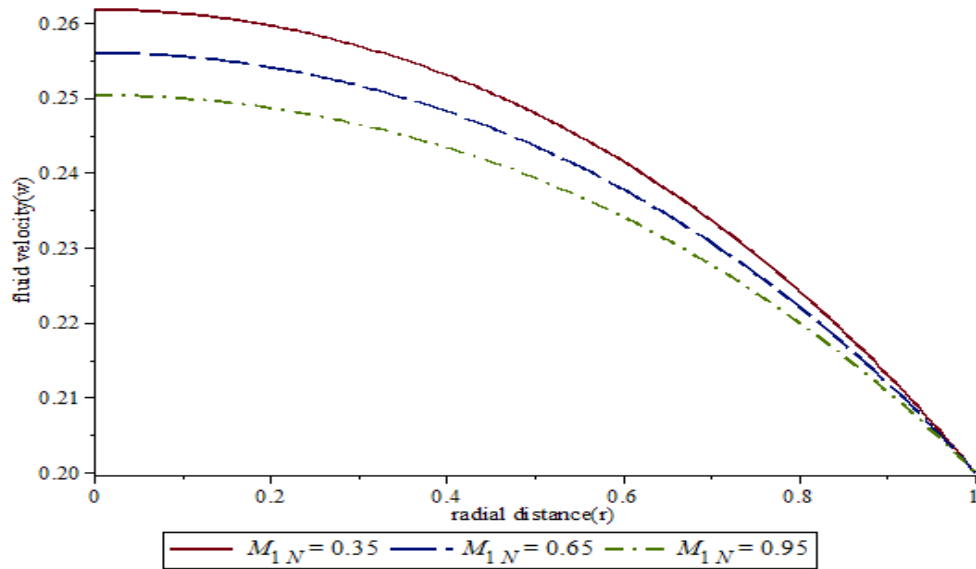


Figure 3b: Variation of Velocity Profile of the Unsteady Blood Flow Model with variable viscosity for various values of the Magnetic Field Parameter in the radial direction.

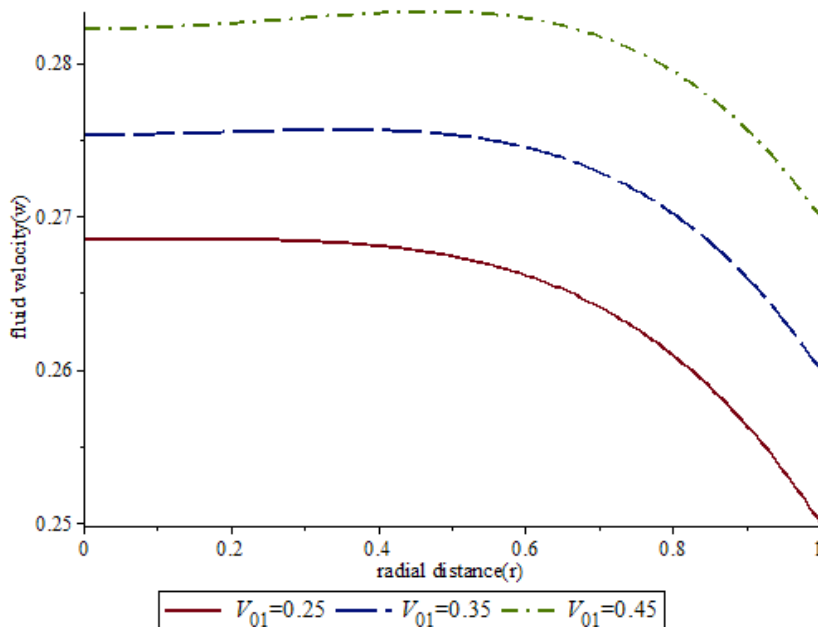


Figure 4a: Variation of Velocity Profile of the Unsteady Blood Flow Model with constant viscosity for various values of the Magnetic Field Parameter in the radial direction.

viscosity for various values of the Slip Velocity in the radial direction.

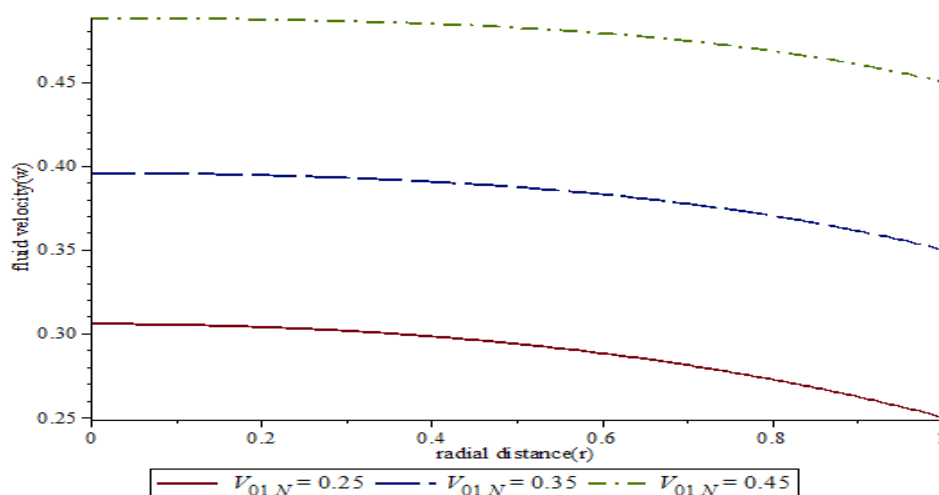


Figure 4b: Variation of Velocity Profile of the Unsteady Blood Flow Model with variable viscosity for various values of the Slip Velocity in the radial direction.

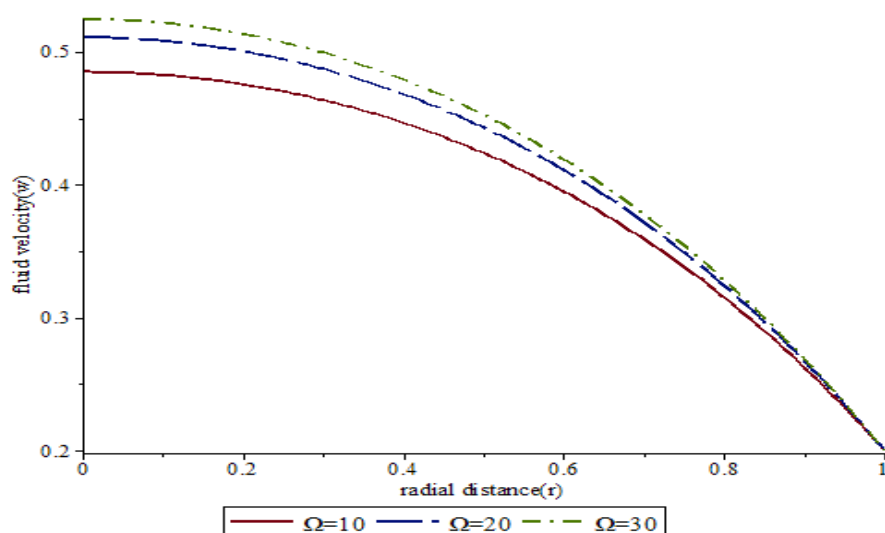


Figure 5a: Variation of Velocity Profile of the Unsteady Blood Flow Model with constant viscosity for various values of the Shear Thinning in the radial direction.

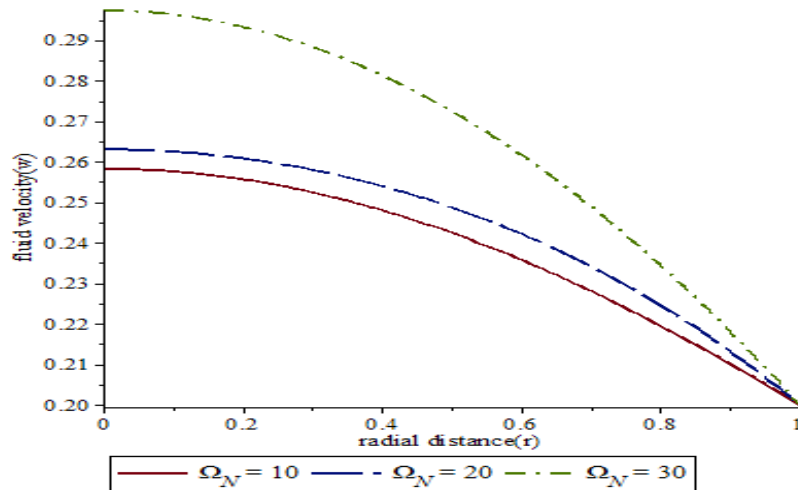


Figure 5b: Variation of Velocity Profile of the Unsteady Blood Flow Model with variable viscosity for various values of the Shear Thinning in the radial direction.

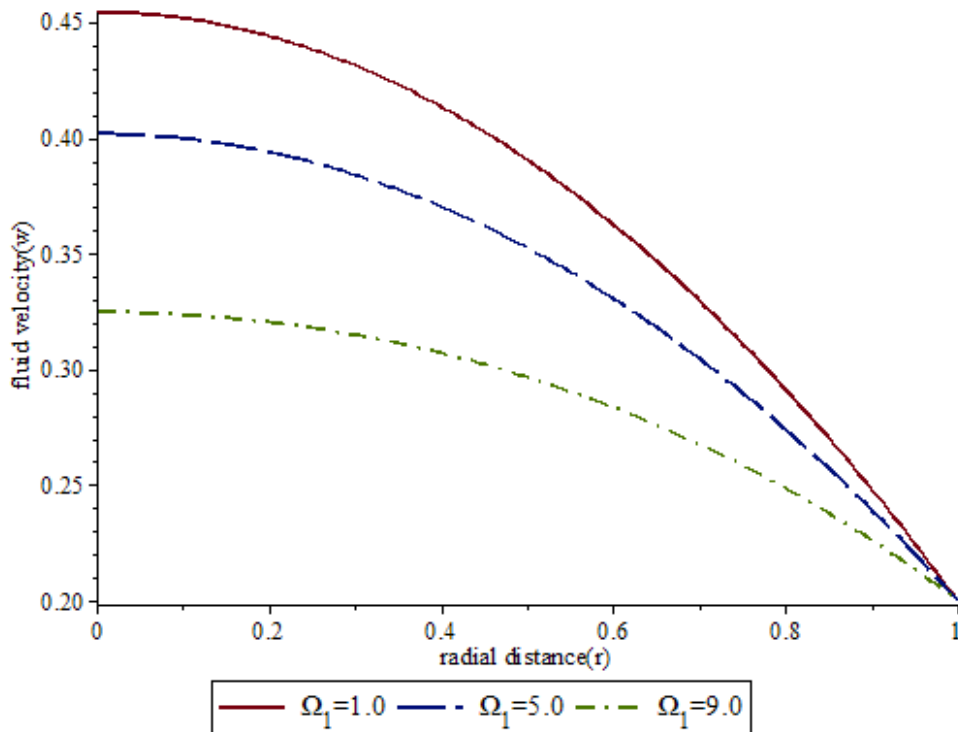


Figure 6a: Variation of Velocity Profile of the Unsteady Blood Flow Model with constant viscosity for various values of the Shear Thinning in the radial direction.

Viscosity for various values of the Shear Thickening in the radial direction.

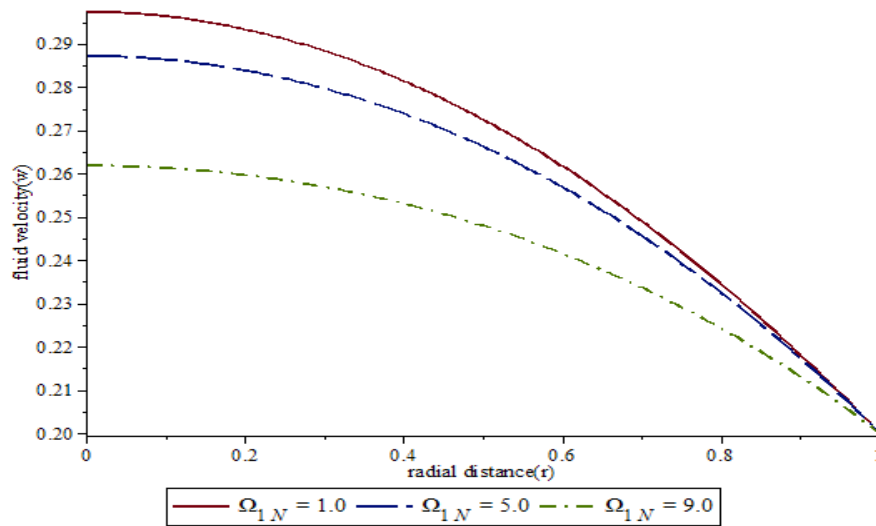


Figure 6b: Variation of Velocity Profile of the Unsteady Blood Flow Model with variable

Viscosity for various values of the Shear Thickening in the radial direction.

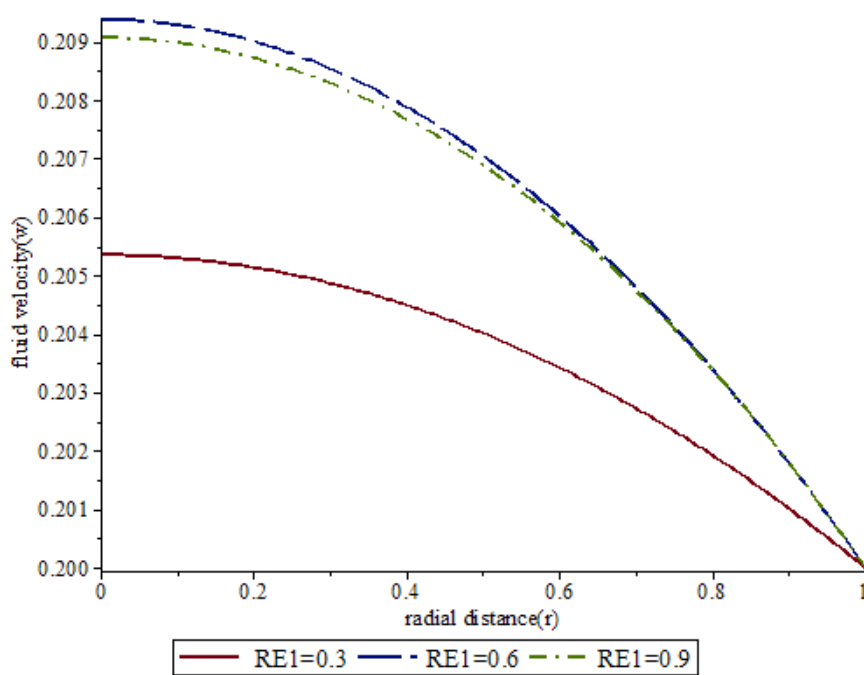


Figure 7a: Variation of Velocity Profile of the Unsteady Blood Flow Model with constant

Viscosity for various values of the Reynold number in the radial direction.

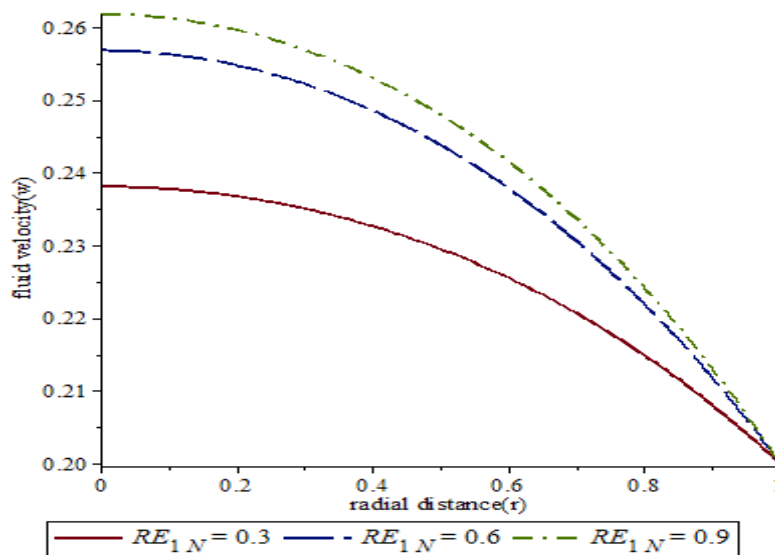


Figure 7b: Variation of Velocity Profile of the Unsteady Blood Flow Model with variable

Viscosity for various values of the Reynold number in the radial direction.

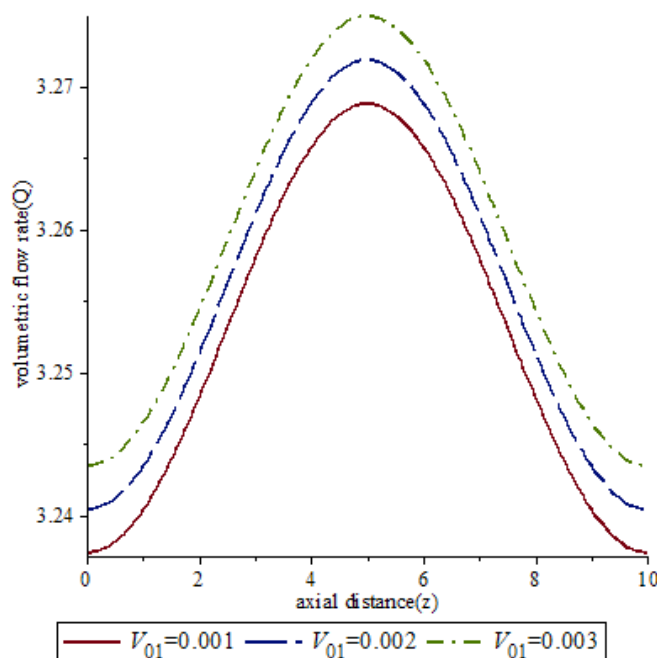


Figure 8a: Variation of Volume Flow Rate of Unsteady Blood Flow Model with constant

Viscosity for the various values of the Slip Velocity in the entire stenotic region along the axial direction.

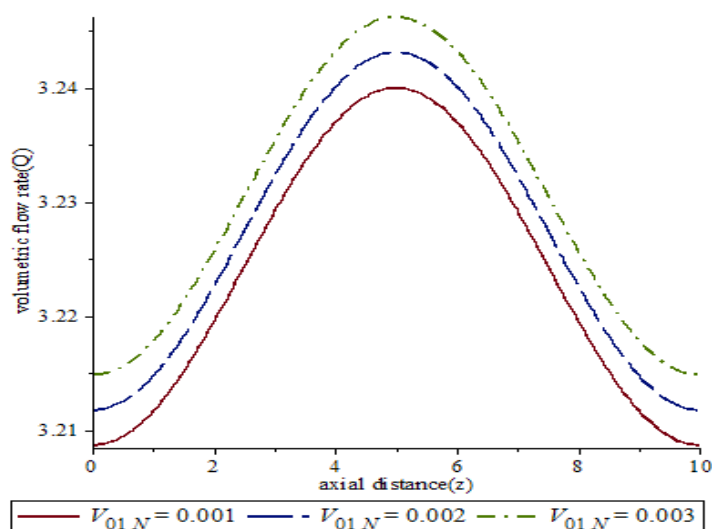


Figure 8b: Variation of Volume Flow Rate of Unsteady Blood Flow Model with variable

Viscosity for the various values of the Slip Velocity in the entire stenotic region along the axial direction.

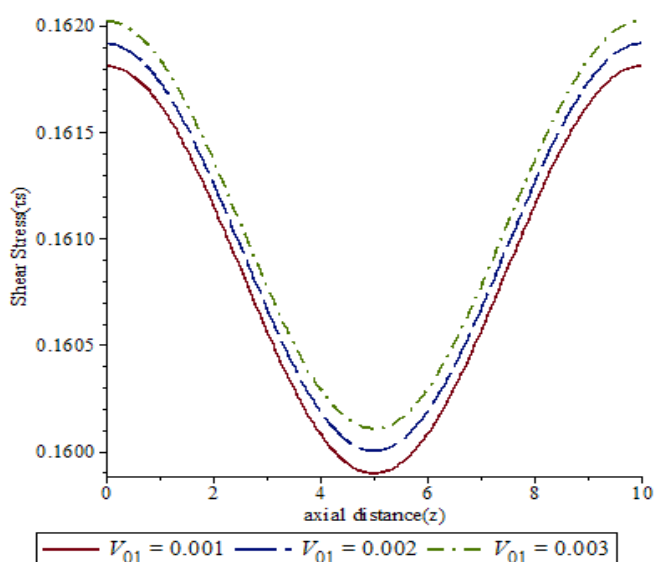


Figure 9a: Variation of Wall Shear Stress of Unsteady Blood Flow Model with constant

Viscosity for the various values of the Slip Velocity in the entire stenotic region along the axial direction.

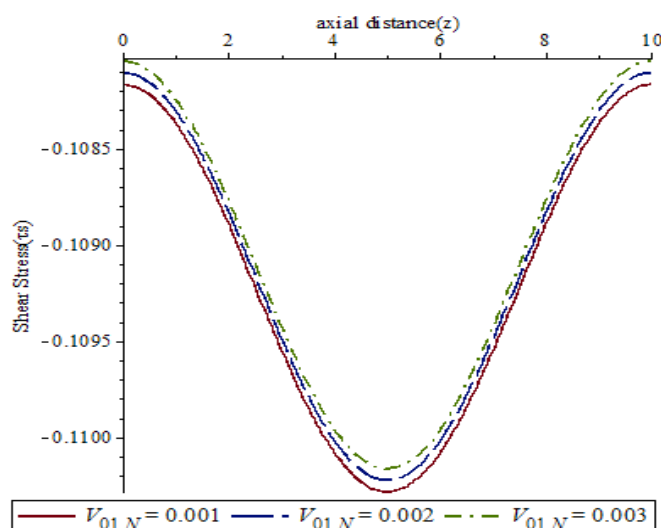


Figure 9b: Variation of Wall Shear Stress of Unsteady Blood Flow Model with variable

Viscosity for the various values of the Slip Velocity in the entire stenotic region along the axial direction.

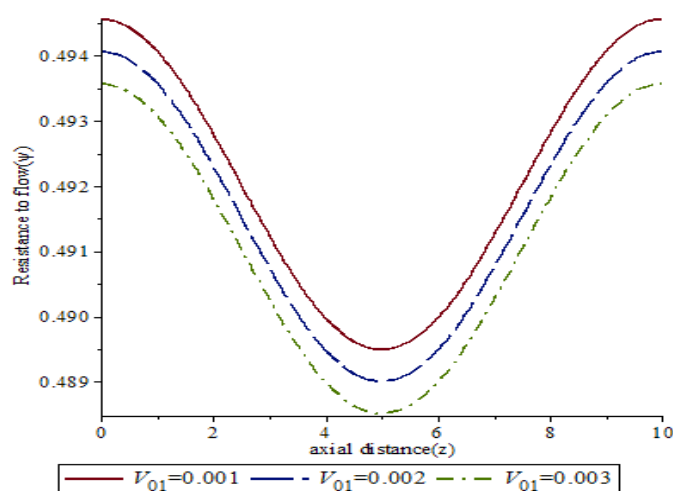


Figure 10a: Variation of Wall Shear Stress of Unsteady Blood Flow Model with constant

Viscosity for various values of the Slip Velocity in the entire stenotic region along the axial direction.

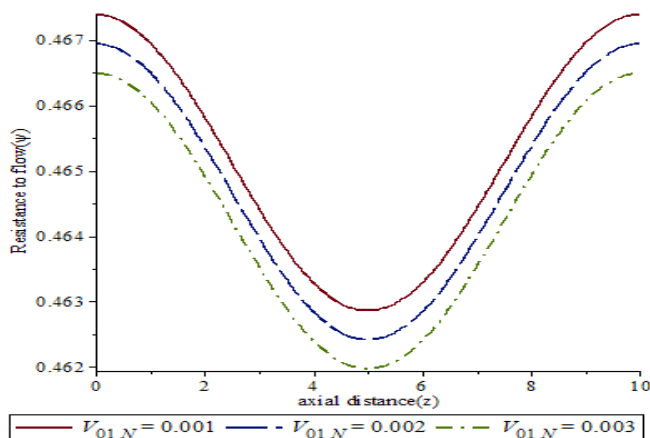


Figure 10b: Variation of Wall Shear Stress of Unsteady Blood Flow Model with variable Viscosity for various values of the Slip Velocity in the entire stenotic region along the axial direction.

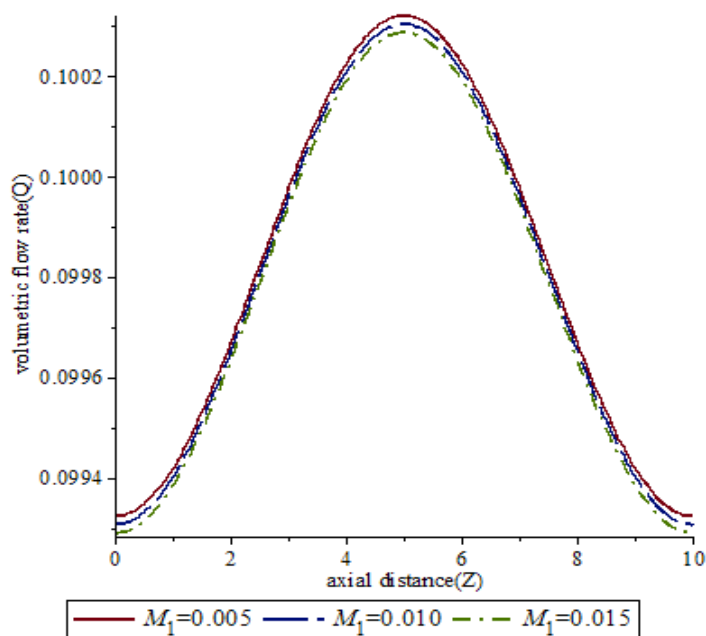


Figure 11a: Variation of Volume Flow Rate of Unsteady Blood Flow Model with constant viscosity for various values of the Magnetic Field Parameter in the entire stenotic region along the axial direction.

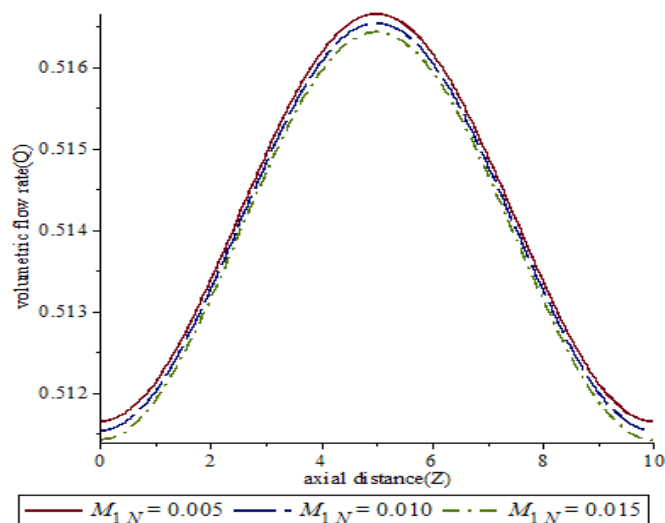


Figure 11a: Variation of Volume Flow Rate of Unsteady Blood Flow Model with constant viscosity for various values of the Magnetic Field Parameter in the entire stenotic region along the axial direction.

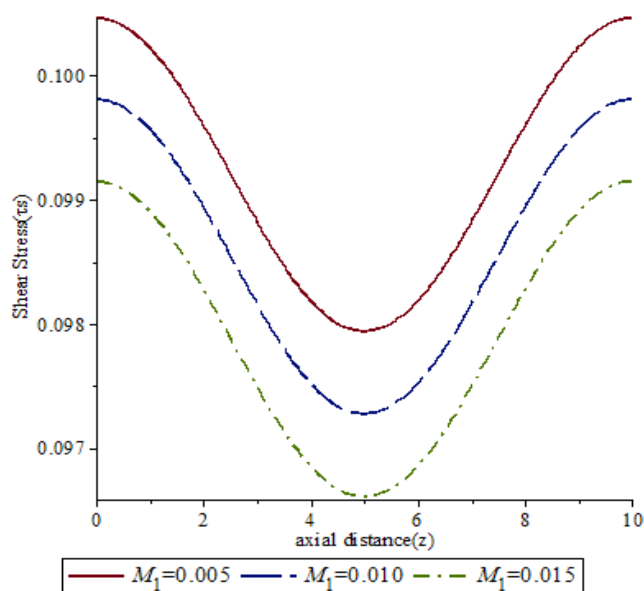


Figure 12a: Variation of Wall Shear Stress of Unsteady Blood Flow Model with constant viscosity for various values of the Magnetic Field Parameter in the entire stenotic region along the axial direction.

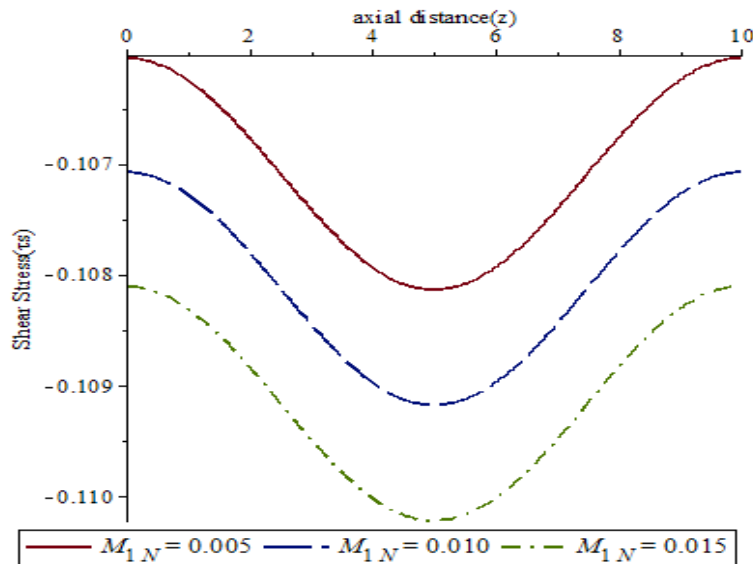


Figure 12b: Variation of Wall Shear Stress of Unsteady Blood Flow Model with variable viscosity for various values of the Magnetic Field Parameter in the entire stenotic region along the axial direction.

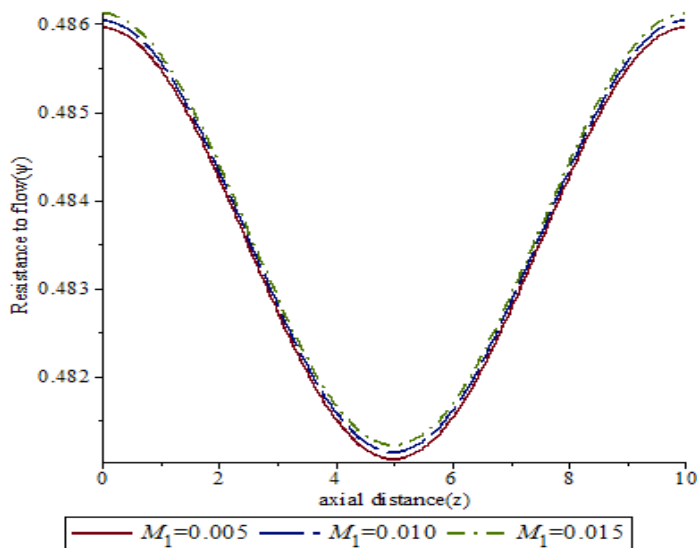


Figure 13a: Variation of Resistance to Unsteady Blood Flow Model with constant viscosity for various values of the Magnetic Field Parameter in the entire stenotic region along the axial direction.

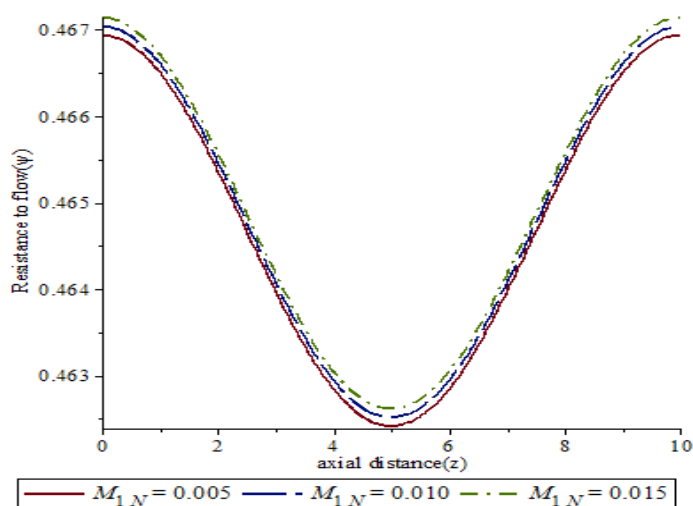


Figure 13b: Variation of Resistance to Unsteady Blood Flow Model with variable viscosity for various values of the Magnetic Field Parameter in the entire stenotic region along the axial direction.

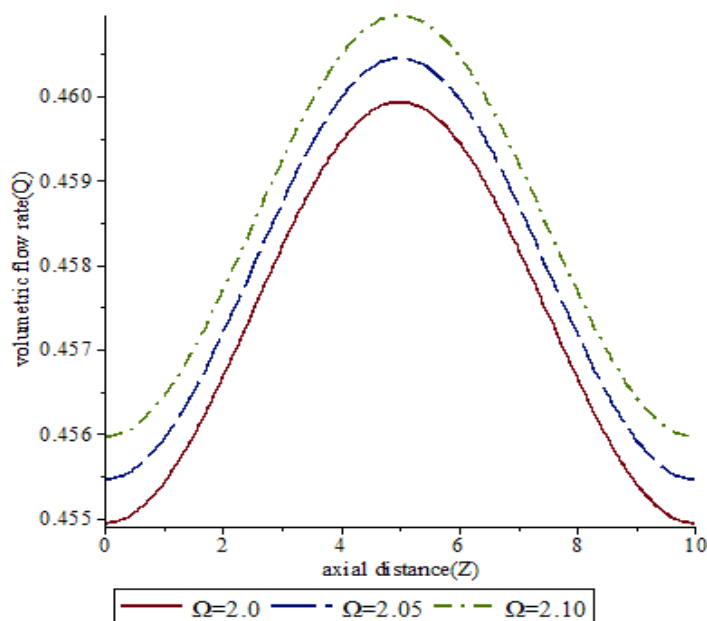


Figure 14a: Variation of Volumetric Flow Rate of Unsteady Blood Flow Model with constant viscosity for various values of the Shear thinning in the entire stenotic region along the axial direction.

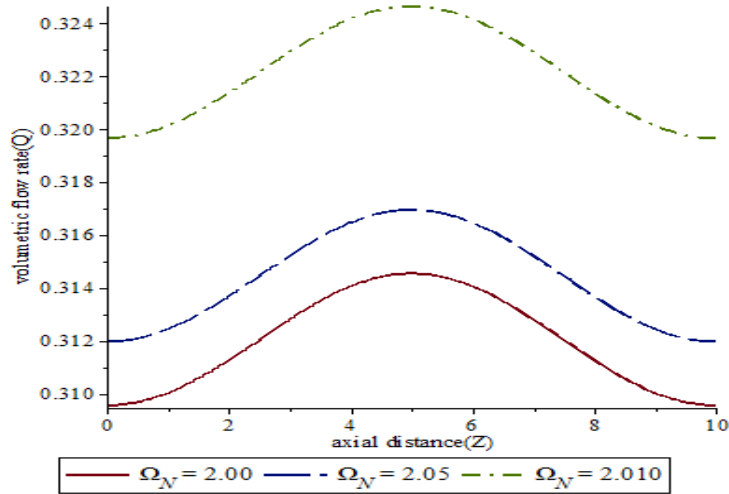


Figure 14b: Variation of Volumetric Flow Rate of Unsteady Blood Flow Model with variable viscosity for various values of the Shear thinning in the entire stenotic region along the axial direction.

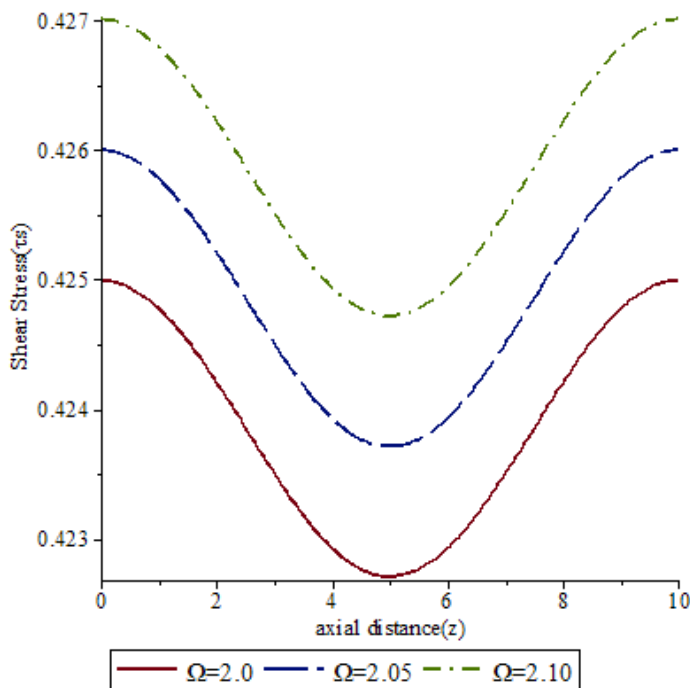


Figure 15a: Variation of Wall Shear Stress of Unsteady Blood Flow Model with constant viscosity for various values of the Shear Thinning in the entire stenotic region along the axial direction.

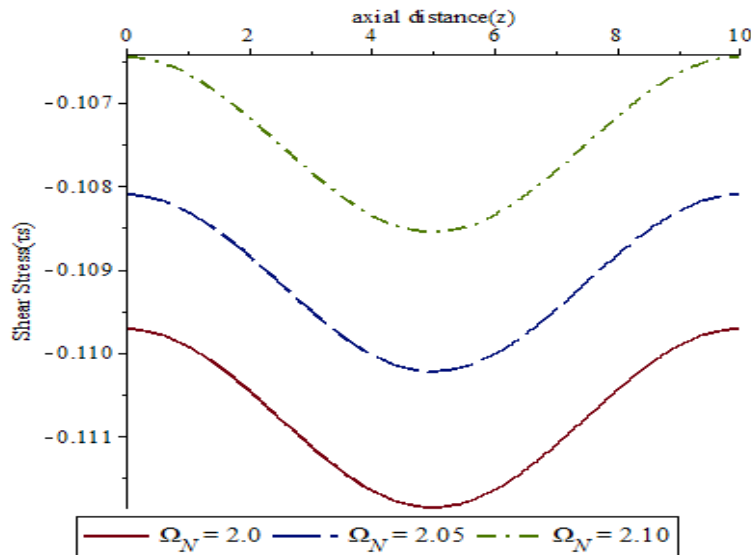


Figure 15b: Variation of Wall Shear Stress of Unsteady Blood Flow Model with variable viscosity for various values of the Shear Thinning in the entire stenotic region along the axial direction.

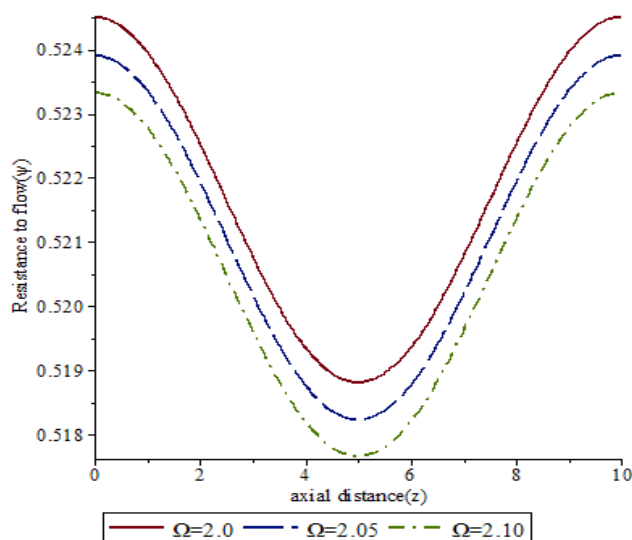


Figure 16a: Variation of Resistance to Unsteady Blood Flow Model with constant viscosity for various values of the Shear Thinning in the entire stenotic region along the axial direction.

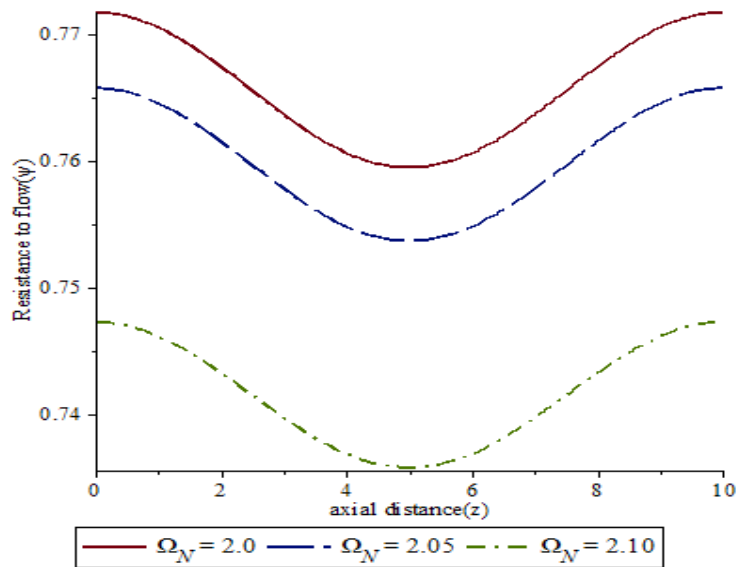


Figure 16b: Variation of Resistance to Unsteady Blood Flow Model with variable viscosity for various values of the Shear Thinning in the entire stenotic region along the axial direction.

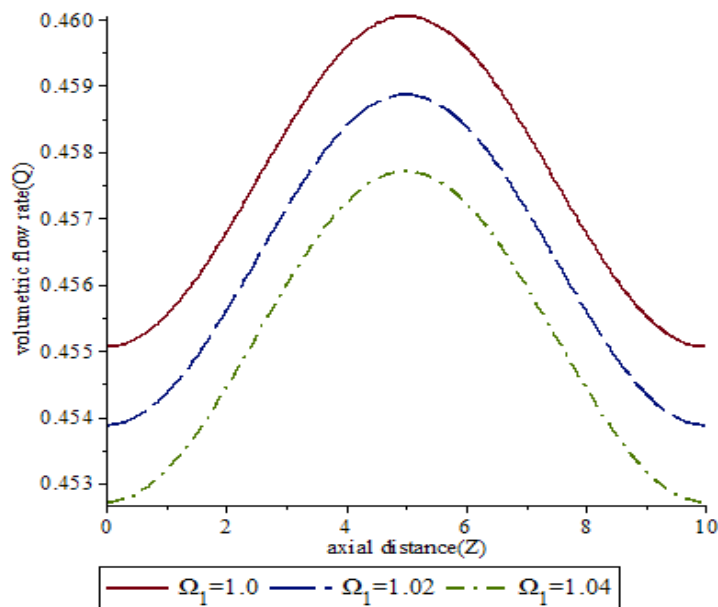


Figure 17a: Variation of Volumetric Flow Rate of Unsteady Blood Flow Model with constant viscosity for various values of the Shear Thickening in the entire stenotic region along the axial direction.

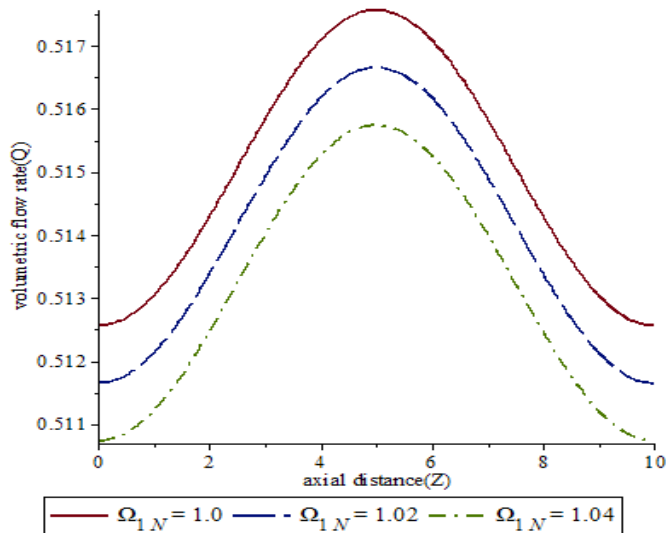


Figure 17b: Variation of Volumetric Flow Rate of Unsteady Blood Flow Model with variable viscosity for various values of the Shear Thickening in the entire stenotic region along the axial direction.

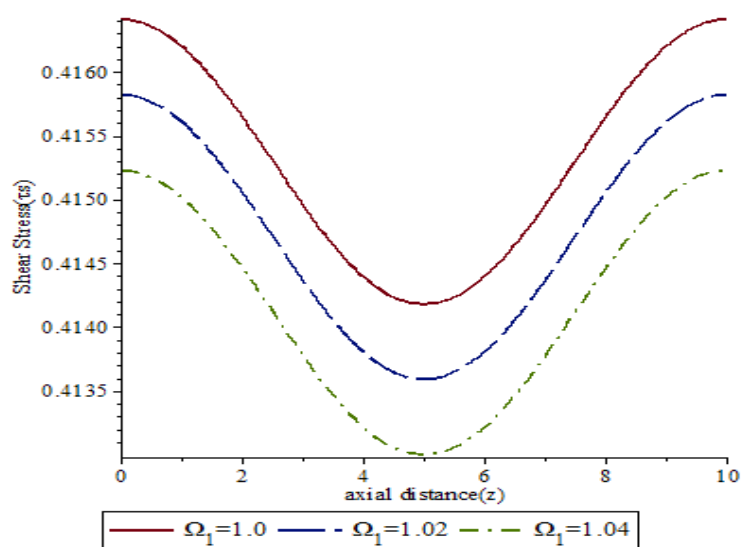


Figure 18a: Variation of Wall Shear Stress of Unsteady Blood Flow Model with constant viscosity for various values of the Shear Thickening in the entire stenotic region along the axial direction.

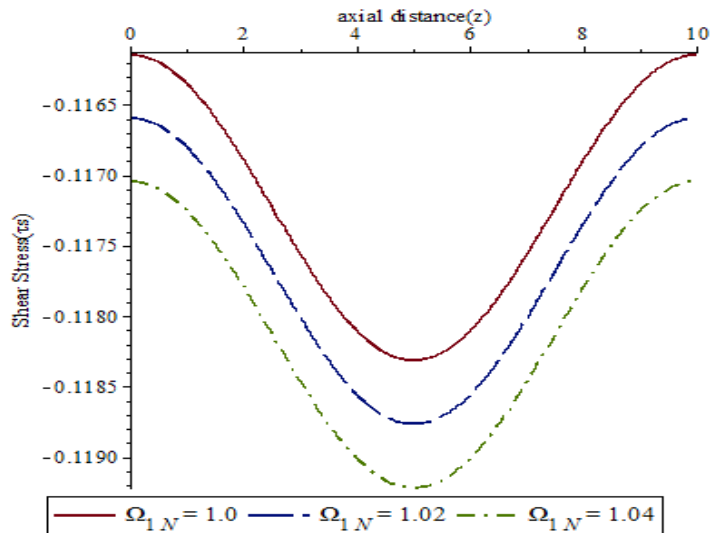


Figure 18b: Variation of Wall Shear Stress of Unsteady Blood Flow Model with variable viscosity for various values of the Shear Thickening in the entire stenotic region along the axial direction.

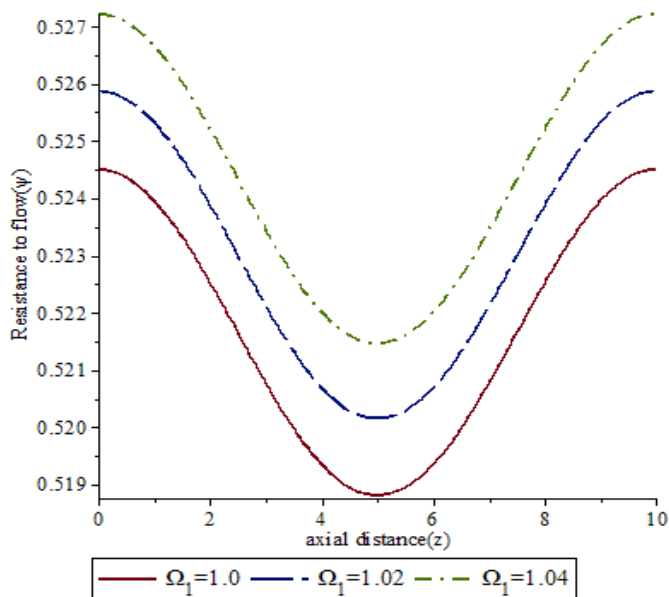


Figure 19a: Variation of Resistance to Unsteady Blood Flow Model with constant viscosity for various values of the Shear Thickening in the entire stenotic region along the axial direction.

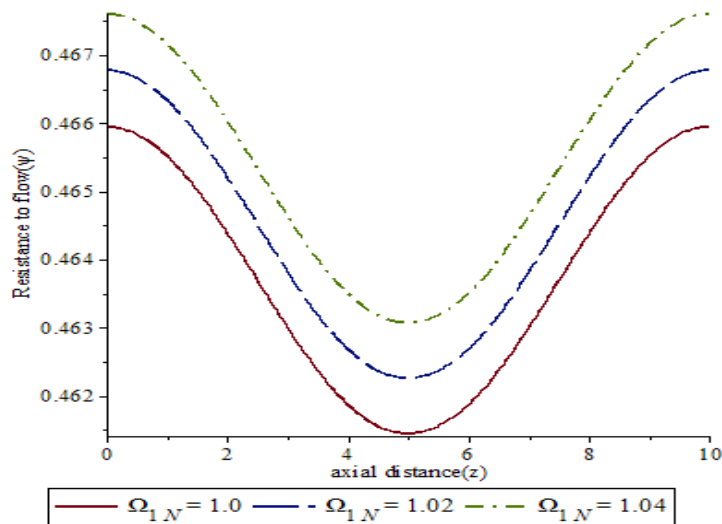


Figure 19b: Variation of Resistance to Unsteady Blood Flow Model with variable viscosity for various values of the Shear Thickening in the entire stenotic region along the axial direction.

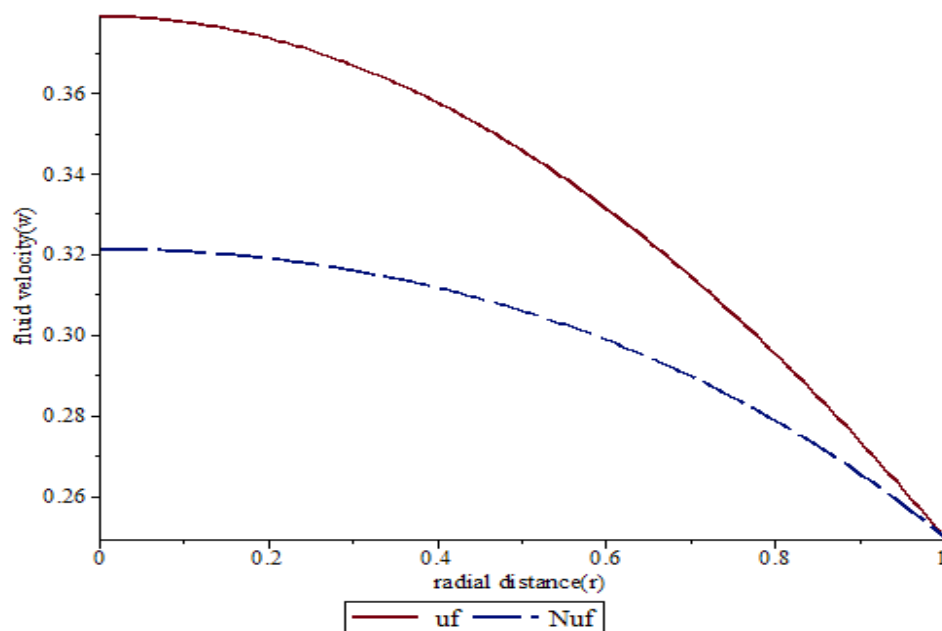


Figure 20: Comparison of the Velocity Profile of Unsteady Blood Flow Models with Constant Viscosity and variable viscosity in the Radial Direction.

RESULTS AND DISCUSSION

In this study, numerical results have been made available to explore the effects of shear thinning, shear thickening, slip velocity and magnetic field parameter on the flow velocity, flow rate, shear stress and flow resistance. The presents mathematical problems are difficult to handle but computation and graphical representations with maple software make it easier. it was reveals that increases in slip velocity and shear thinning significantly lead to an increases in flow velocity, flow rate and shear stress but decrease the resistance to fluid flow for both flow model with constant viscosity and variable viscosity and these are shown in figures 4a, 4b, 8a, 8b, 9a, 9b, 10a,10b, 5a, 5b, 14a, 14b, 15a, 15b, 16a and 16b respectively. Increases in magnetic field parameter and shear thickening lead to deceases in flow velocity, flow rate, and shear stress but increase the resistances to fluid flow for the flow model with constant viscosity and variable viscosity and these are shown in figures 3a, 3b, 11a, 11b, 12a, 12b, 13a,13b, 6a, 6b, 17a, 17b, 18a, 18b, 19a and 19b respectively. Pressure gradient and Reynold number remarkably increases with flow velocity as shown in figures 2a, 2b, 7a and 7b respectively. Finally, the flow velocity of unsteady blood flow model with constant viscosity is higher than that with variable viscosity as indicated in figure 20.

Conclusion

A theoretical study of unsteady blood flow with constant and variable viscosities through a stenosed artery using a third grade fluid model has been carried out. In this study slip velocity and externally magnetic field are taken into consideration. The problems are solved using Galerkin weighted residual and forth order Runge-Kutta methods. The main findings of the present study may be listed as follows:

- (i) The flow velocity, flow rate and shear stress decreases while flow resistance decreases with the increases of magnetic field strength and shear thickening for the blood flow models with constant and variable viscosities.
- (ii) The flow velocity, flow rate and shear stress increases while flow resistance decreases with the increases of slip velocity and shear thinning for the blood flow models with constant and variable viscosities.
- (iii) The effects of shear thinning, slip velocity, magnetic field and pressure gradient on the flow velocity are more pronounce and remarkable for the unsteady blood flow model with variable viscosity when compare to unsteady blood flow model with constant viscosity.
- (iv) Finally, with fixed values of all the parameters, the velocity profile of the unsteady blood flow model with constant viscosity is higher than that with variable viscosity.

The research analysis incorporating slip velocity in the constricted artery can help to reduce blood viscosity because high blood viscosity is a risk factor in the cardiovascular disorder. Also, incorporating externally applied magnetic field will be useful for the reduction of blood flow during surgery and magnetic resonance imaging (MRI).

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Nomenclatures

w - Fluid velocity	\bar{w} - Dimensionless fluid velocity
t - Time component	\bar{t} - Dimensionless time component
r - Radial distance	y - Dimensionless radial distance
z - Axial distance	w_s - Slip velocity
V_{01} - Dimensionless slip velocity for the unsteady blood flow with constant viscosity	
V_{01N} - Dimensionless slip velocity for the unsteady blood flow with variable viscosity	
R_0 - Radius of the normal artery	β_0 - Magnetic Field Strength
$R(z)$ - Radius of the artery in a stenotic region	τ_s - Wall Shear Stress
ψ - Resistance to flow	ξ - Maximum height of the stenosis
Q - Volumetric flow rate	L - Length of the stenosis
W = Fluid velocity	
G_1 -Pressure gradient for the unsteady flow with constant viscosity	
G_{1N} -Pressure gradient for the unsteady flow with variable viscosity	
M_1 -Magnetic field parameter for the unsteady flow with constant viscosity	
M_{1N} -Magnetic field parameter for the unsteady flow with variable viscosity	
Ω -Shear thinning for the unsteady flow with constant viscosity	
Ω_N -Shear thinning for the unsteady flow with variable viscosity	
Ω_1 -Shear thickening for the unsteady flow with constant viscosity	
Ω_{1N} -Shear thickening for the unsteady flow with variable viscosity	
RE_1 —Reynold number for the unsteady flow with constant viscosity	
RE_{1N} —Reynold number for the unsteady flow with variable viscosity	

